

# Deformed nuclear halos

Shan-Gui Zhou (周善贵)

Institute of Theoretical Physics, Chinese Academy of Sciences, Beijing

School of Physical Sciences, University of Chinese Academy of Sciences, Beijing

Center of Theoretical Nucl. Phys., National Laboratory of Heavy Ion Accelerator, Lanzhou

Synergetic Innovation Center for Quantum Effects & Application, Hunan Normal Univ., Changsha

*Collaborators:* Lulu Li (ITP, PKU, IAPCM), Jie Meng (PKU),  
P. Ring (TU Munich & PKU), Xiang-Xiang Sun (ITP)  
Jie Zhao (ITP), En-Guang Zhao (ITP)

*Supported by:* NSFC, CAS & MOST;  
HPC Cluster of KLTP/ITP-CAS  
ScGrid of CNIC-CAS

# Happy birthday, Professor Arima!



Prof. Arima's visit of  
ITP/CAS in 1998

Congratulations also from  
Prof. En-Guang Zhao  
(赵恩广)



# What I learned from Prof. Arima (continued)

The screenshot shows the Beihang University News website (Beihang University Website) with the following details:

- Header:** 北京航空航天大学 BEIHANG UNIVERSITY 新闻网
- Top navigation:** 北航官网 旧版新闻网 ENGLISH iHome
- Secondary navigation:** 首页 综合新闻 专题新闻 校园风采 科教在线 媒体北航 光影北航 视频新闻 文艺园地 信息公告 学术及文化活动 网上投稿
- Breadcrumb:** / 学术及文化活动
- Section title:** 【北航大讲堂预告】有马朗人教授将莅校做“自然、文化和科学中的对称性”专题讲座
- Print button:** 打印

## Symmetries in Arts, Culture and Nature

Akito Arima

# What I learned from Prof. Arima (continued)

## Summary of this part

- 1 European paintings and buildings do not break Bilateral Symmetry
- 2 Chinese buildings keep bilateral symmetry but Chinese paintings break it
- 3 Japanese buildings and paintings often break bilateral symmetry

有马先生：北航大讲堂  
2010年4月28日

## A Question ; Why?

One answer is because culture is affected by the climate such as monsoon in Asia, dry air in the Middle East and African desert land and moderate climate in Europe.

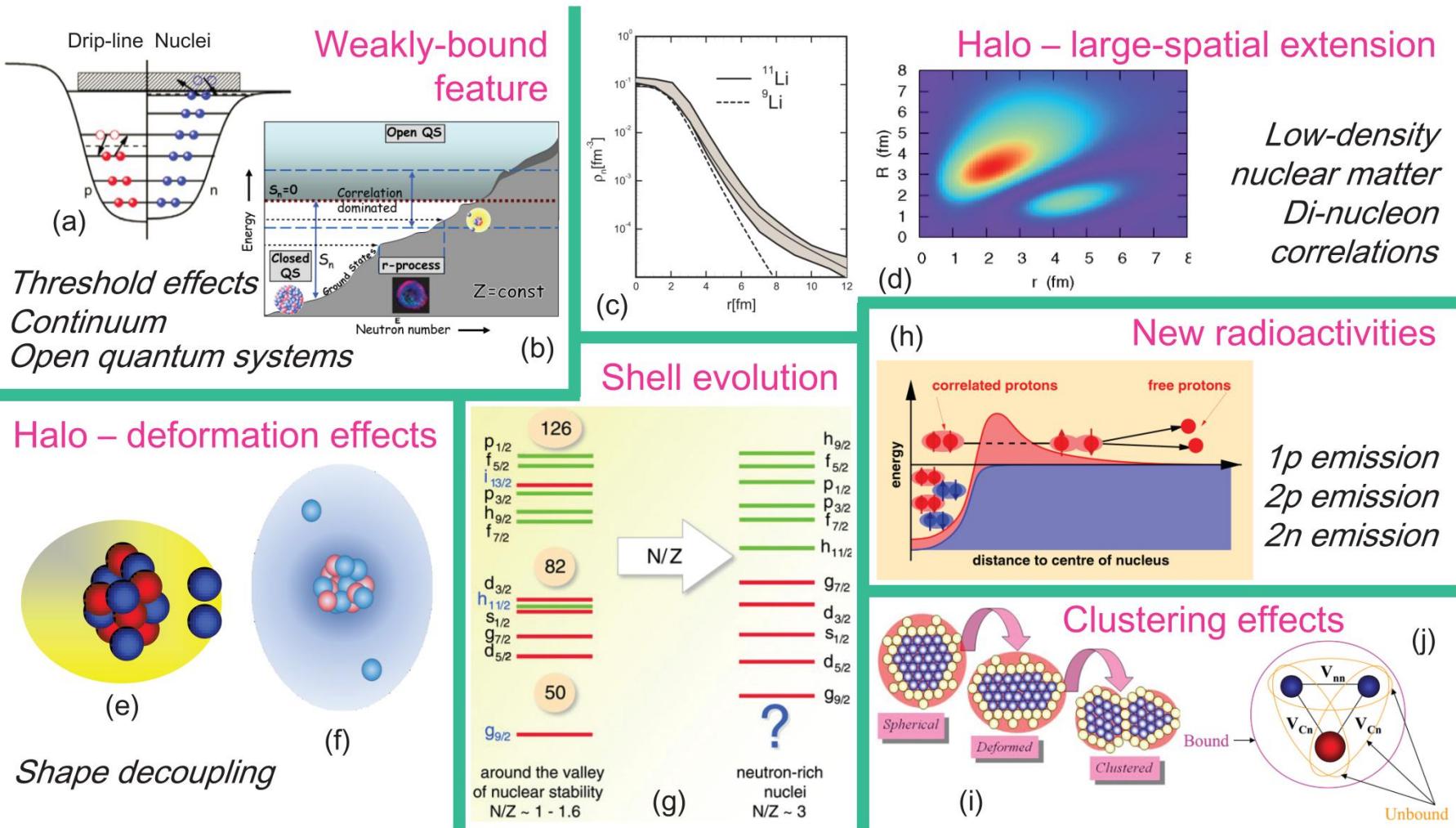
One more question. Can such differences affect our thoughts such as physics?

# Contents

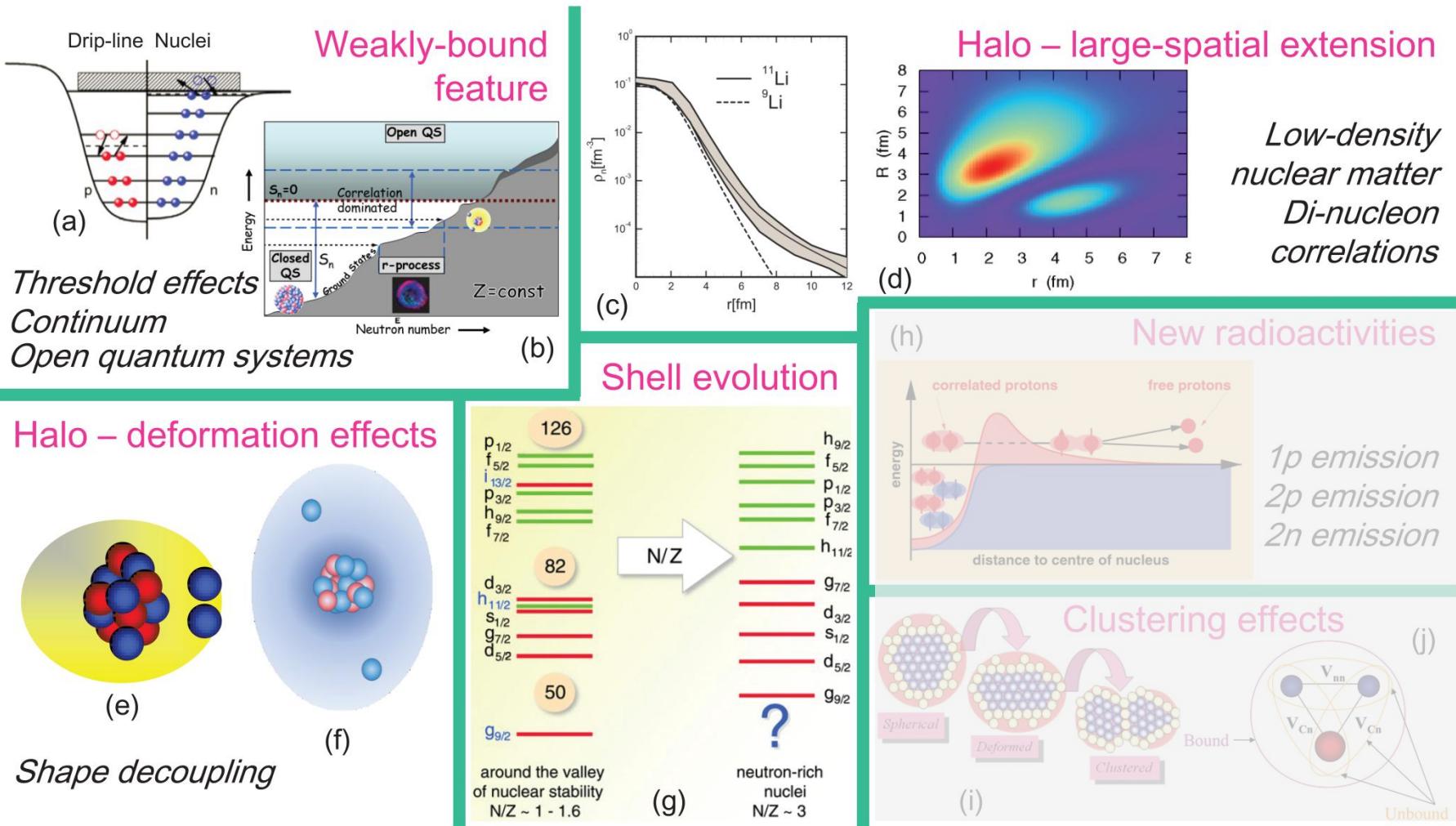
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- Introduction
- Deformed RHB model in continuum (Woods-Saxon basis)
- Shape decoupling in deformed halo nuclei
  - $^{44}\text{Mg}$ : prolate core but oblate halo
  - $^{22}\text{C}$ : oblate core but prolate halo
  - $^{11}\text{Li}$ ,  $^{22}\text{C}$  &  $^{44}\text{Mg}$ : triangle of Borromean nuclei
- How to probe shape decoupling in deformed halo nuclei?
- Summary & perspectives

# Physics in exotic nuclear structure

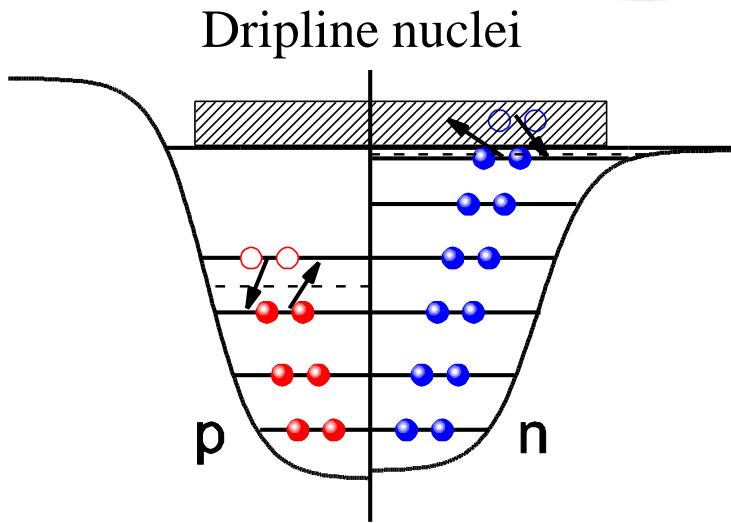
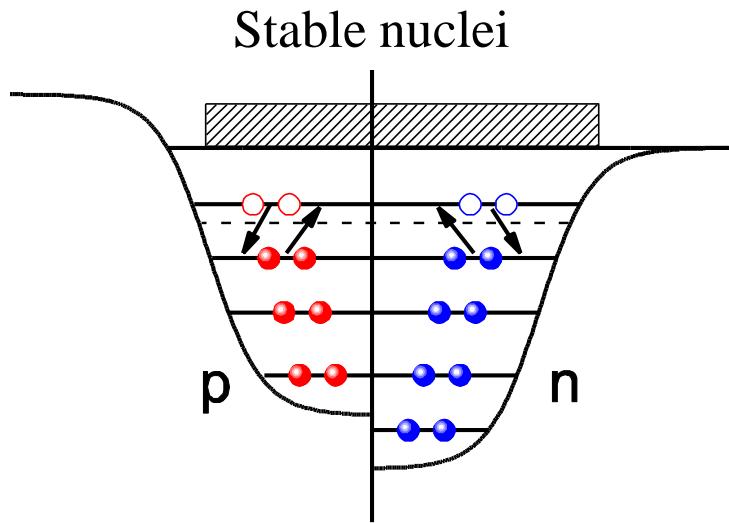
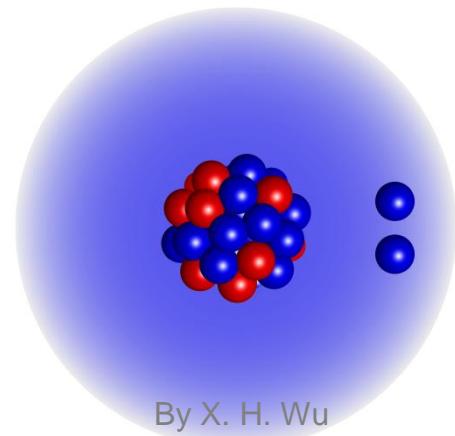


# Physics in exotic nuclear structure



# Characteristics of halo nuclei

- Weakly bound; large spatial extension
- Continuum can not be ignored



## *Self-consistent description:*

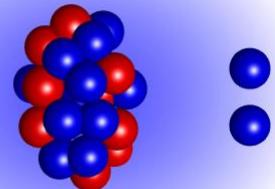
- Weakly bound, continuum
- Large spatial distribution
- Couplings among ...

Meng\_Toki\_SGZ\_Zhang\_Long\_Geng2006  
Prog. Part. Nucl. Phys. 57-470  
Meng & SGZ 2015, J. Phys. G42-093101

Bulgac1980; nucl-th/9907088

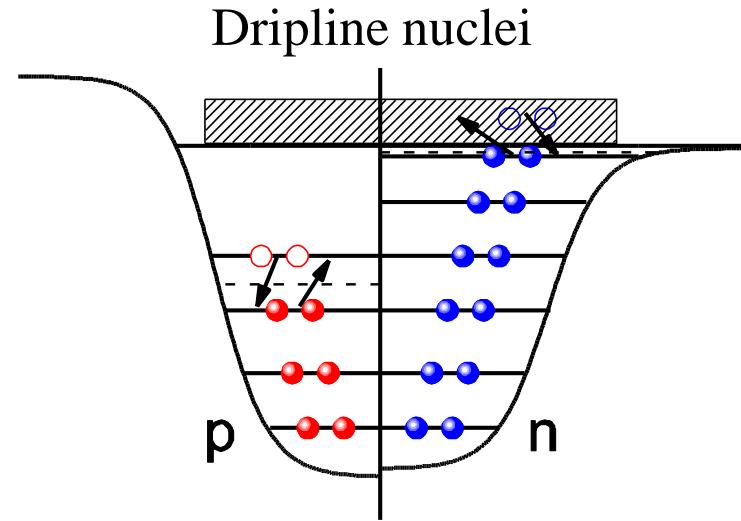
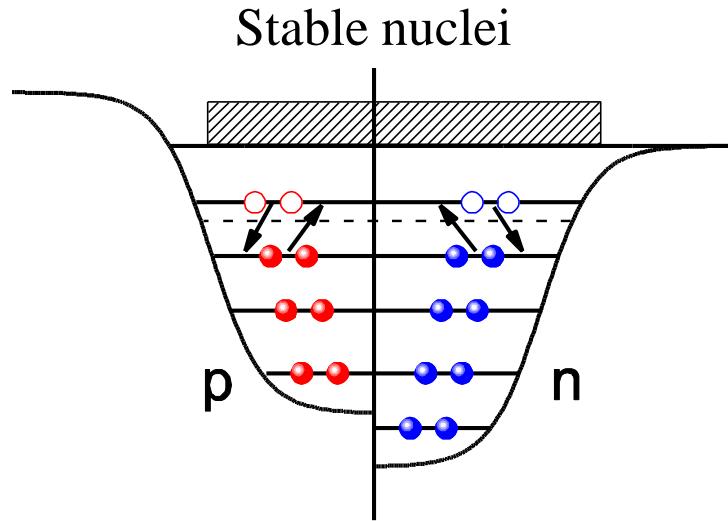
Dobaczewski\_Flocard\_Treiner1984\_NPA422-103

# Characteristics of deformed halo nuclei



By X. H. Wu

- Weakly bound; large spatial extension
- Continuum can not be ignored



## *Self-consistent description:*

- Weakly bound, continuum
- Large spatial distribution
- Deformation effects**
- Couplings among ...

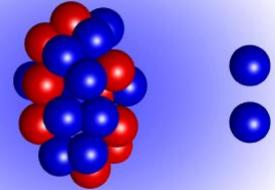
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Bulgac1980; nucl-th/9907088

Dobaczewski\_Flocard\_Treiner1984\_NPA422-103

# What we aim at

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By X. H. Wu

*A self-consistent description of*

- ✓ Deformation
- ✓ Continuum contribution
- ✓ Large spatial distribution
- ✓ Interplays among them

by developing a  
relativistic Hartree-Bogoliubov model

# Covariant Density Functional Theory (CDFT)

---

$$\begin{aligned}
\mathcal{L} = & \bar{\psi}_i (i\cancel{\partial} - M) \psi_i + \frac{1}{2} \partial_\mu \sigma \partial^\mu \sigma - U(\sigma) - g_\sigma \bar{\psi}_i \sigma \psi_i \\
& - \frac{1}{4} \Omega_{\mu\nu} \Omega^{\mu\nu} + \frac{1}{2} m_\omega^2 \omega_\mu \omega^\mu - g_\omega \bar{\psi}_i \cancel{\omega} \psi_i \\
& - \frac{1}{4} \vec{R}_{\mu\nu} \vec{R}^{\mu\nu} + \frac{1}{2} m_\rho^2 \vec{\rho}_\mu \vec{\rho}^\mu - g_\rho \bar{\psi}_i \cancel{\rho} \vec{\tau} \psi_i \\
& - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - e \bar{\psi}_i \frac{1 - \tau_3}{2} \cancel{A} \psi_i,
\end{aligned}$$

Serot\_Walecka1986\_ANP16-1

Reinhard1989\_RPP52-439

Ring1996\_PPNP37-193

Vretenar\_Afanasjev\_Lalazissis\_Ring2005\_PR409-101

Meng\_Toki\_SGZ\_Zhang\_Long\_Geng2006\_PPNP57-470

$$(\alpha \cdot \mathbf{p} + \beta(M + S(\mathbf{r})) + V(\mathbf{r})) \psi_i = \epsilon_i \psi_i$$

Liang\_Meng\_SGZ2015\_PR570-1

$$(-\nabla^2 + m_\sigma^2) \sigma = -g_\sigma \rho_S - g_2 \sigma^2 - g_3 \sigma^3$$

Meng\_SGZ2015\_JPG42-093101

$$(-\nabla^2 + m_\omega^2) \omega = g_\omega \rho_V - c_3 \omega^3$$

$$(-\nabla^2 + m_\rho^2) \rho = g_\rho \rho_3$$

$$-\nabla^2 A = e \rho_C$$

# Deformed RHB theory in continuum

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$$\sum_{\sigma' p'} \int d^3 r' \begin{pmatrix} h_D(\mathbf{r}\sigma p, \mathbf{r}\sigma' p') - \lambda & \Delta(\mathbf{r}\sigma p, \mathbf{r}'\sigma' p') \\ -\Delta^*(\mathbf{r}\sigma p, \mathbf{r}'\sigma' p') & -h_D(\mathbf{r}\sigma p, \mathbf{r}\sigma' p') + \lambda \end{pmatrix} \begin{pmatrix} U_k(\mathbf{r}'\sigma' p') \\ V_k(\mathbf{r}'\sigma' p') \end{pmatrix} = E_k \begin{pmatrix} U_k(\mathbf{r}\sigma p) \\ V_k(\mathbf{r}\sigma p) \end{pmatrix}$$

Kucharek\_Ring1991\_ZPA339-23

$$U_k(\mathbf{r}\sigma p) = \sum_{i\kappa} \begin{pmatrix} u_{k,(i\kappa)}^{(m)} \varphi_{i\kappa m}(\mathbf{r}\sigma p) \\ u_{k,(i\tilde{\kappa})}^{(\bar{m})} \tilde{\varphi}_{i\kappa m}(\mathbf{r}\sigma p) \end{pmatrix}$$

Woods-Saxon basis

$$V_k(\mathbf{r}\sigma p) = \sum_{i\kappa} \begin{pmatrix} v_{k,(i\kappa)}^{(m)} \varphi_{i\kappa m}(\mathbf{r}\sigma p) \\ v_{k,(i\tilde{\kappa})}^{(\bar{m})} \tilde{\varphi}_{i\kappa m}(\mathbf{r}\sigma p) \end{pmatrix}$$

$$\varphi_{i\kappa m}(\mathbf{r}\sigma) = \frac{1}{r} \begin{pmatrix} iG_{i\kappa}(r)Y_{jm}^l(\Omega\sigma) \\ -F_{i\kappa}(r)Y_{jm}^{\tilde{l}}(\Omega\sigma) \end{pmatrix}$$

$$\begin{pmatrix} \mathcal{A} & \mathcal{B} \\ \mathcal{C} & \mathcal{D} \end{pmatrix} \begin{pmatrix} \mathcal{U} \\ \mathcal{V} \end{pmatrix} = E \begin{pmatrix} \mathcal{U} \\ \mathcal{V} \end{pmatrix}$$

$$\mathcal{U} = \left( u_{k,(i\kappa)}^{(m)} \right), \quad \mathcal{V} = \left( v_{k,(i\tilde{\kappa})}^{(\bar{m})} \right)$$

SGZ\_Meng\_Ring 2007\_ISPUN Proc.

SGZ\_Meng\_Ring\_Zhao 2010\_PRC82-011301R

SGZ\_Meng\_Ring\_Zhao 2011\_JPConfProc312-092067

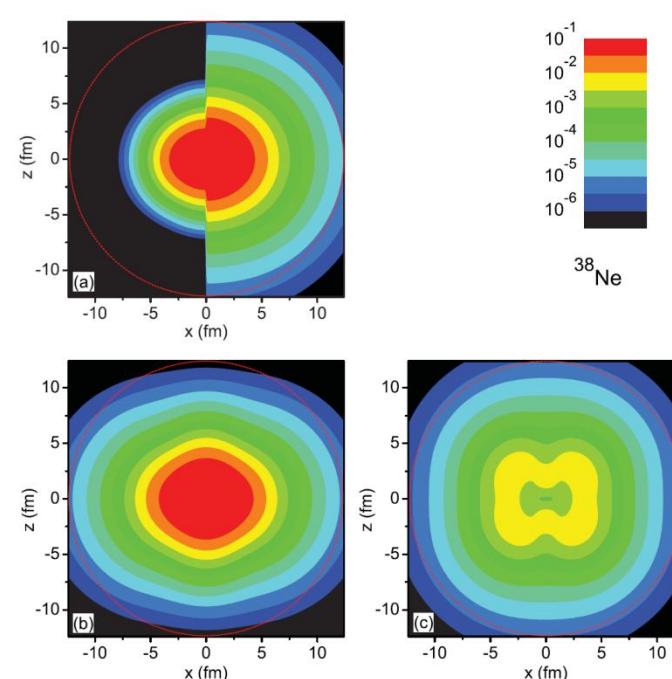
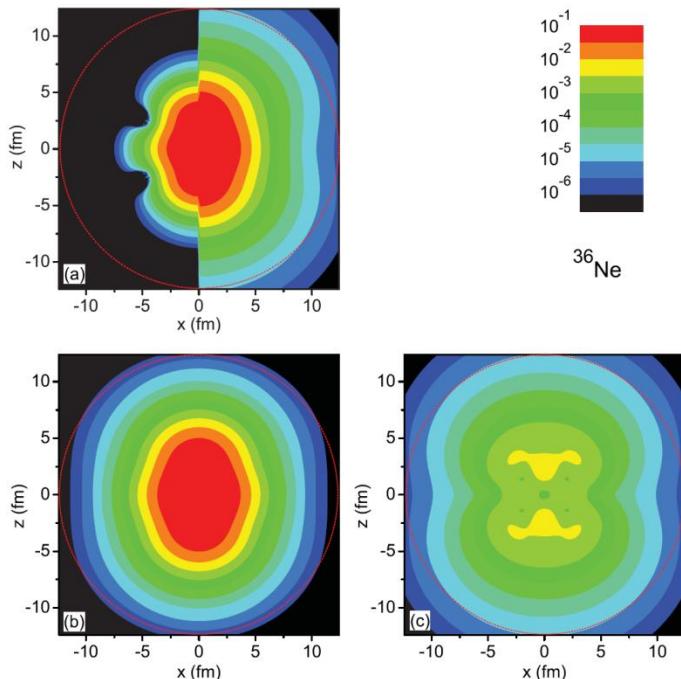
Li\_Meng\_Ring\_Zhao\_SGZ 2012\_PRC85-024312

Li\_Meng\_Ring\_Zhao\_SGZ 2012\_ChinPhysLett29-042101

# Conditions for occurrence of a halo & its shape

- Existence & deformation of neutron halo depend on quantum numbers of the main components of the s.p. orbits around Fermi surface
  - s levels with  $\Lambda = 0 \Rightarrow$  spherical halos
  - p levels with  $\Lambda = 0 \Rightarrow$  prolate halos
  - p levels with  $\Lambda = 1 \Rightarrow$  oblate halos
  - d, f, ... levels: no halos

SGZ\_Meng\_Ring\_Zhao 2010  
PRC82-011301R  
Li\_Meng\_Ring\_Zhao\_SGZ 2012  
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[SGZ\\_Meng\\_Ring\\_Zhao 2010](#)

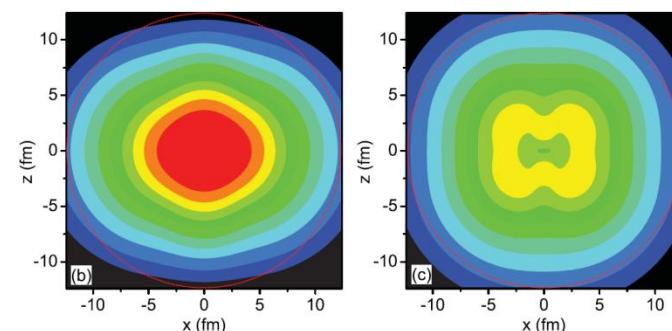
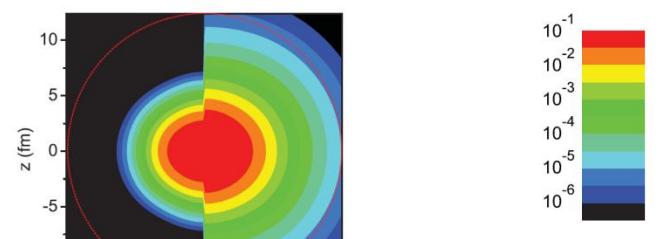
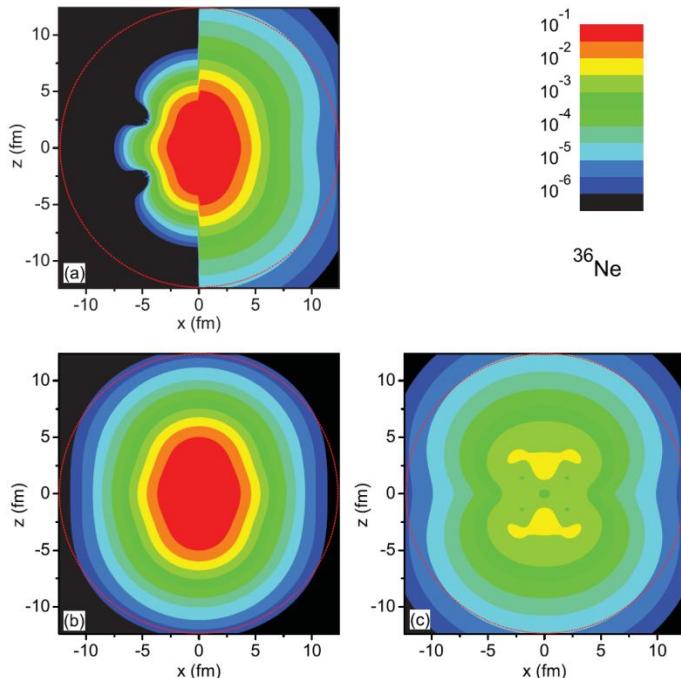
PRC82-011301R

[Li\\_Meng\\_Ring\\_Zhao\\_SGZ 2012](#)

PRC85-024312

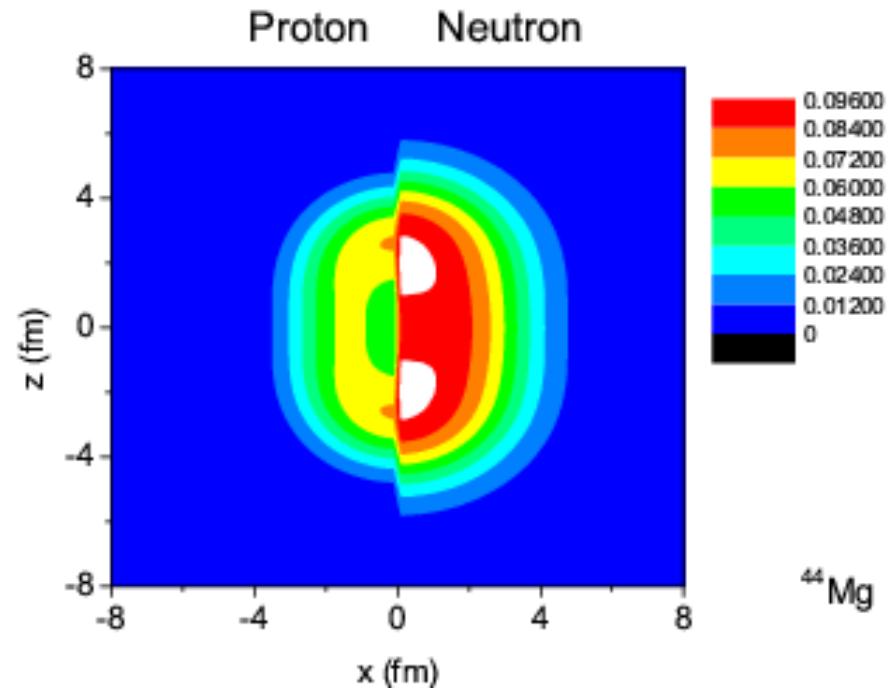
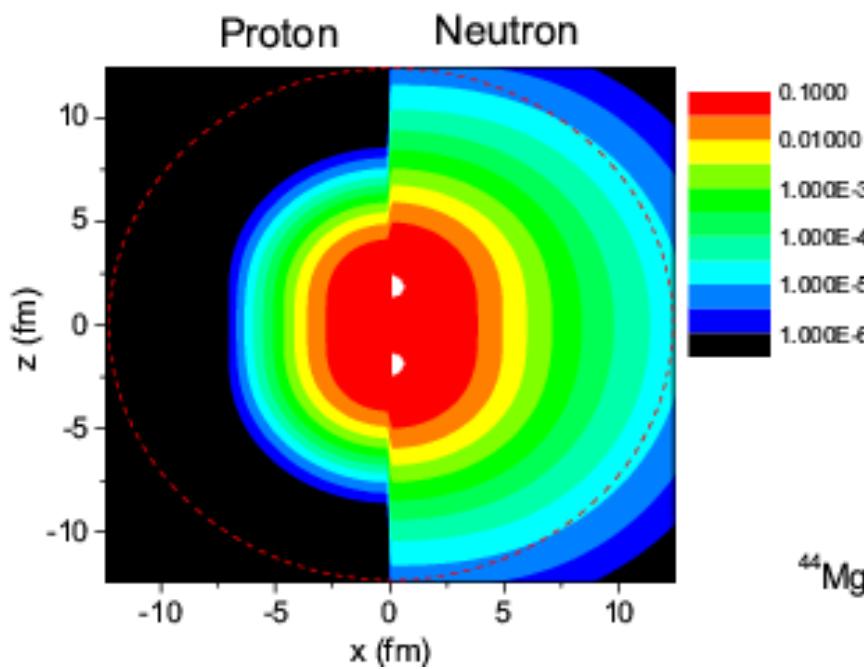
[Pei\\_Zhang\\_Xu2013PRC87-051302R](#)

[Nakada\\_Takayama2018\\_PRC98-011301R](#)



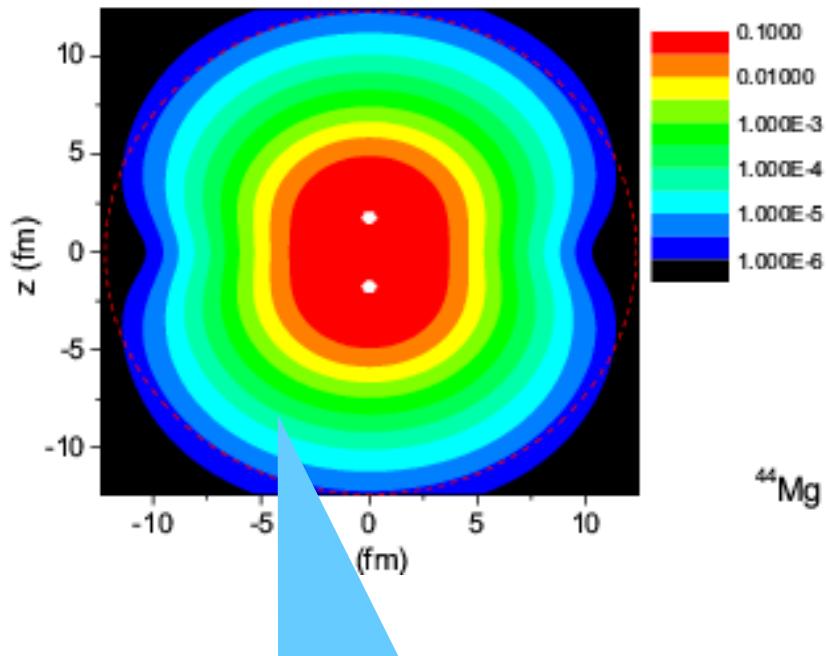
# $^{44}\text{Mg}$ : Density distributions

SGZ\_Meng\_Ring\_Zhao 2010 PRC82-011301R  
Li\_Meng\_Ring\_Zhao\_SGZ 2012 PRC85-024312

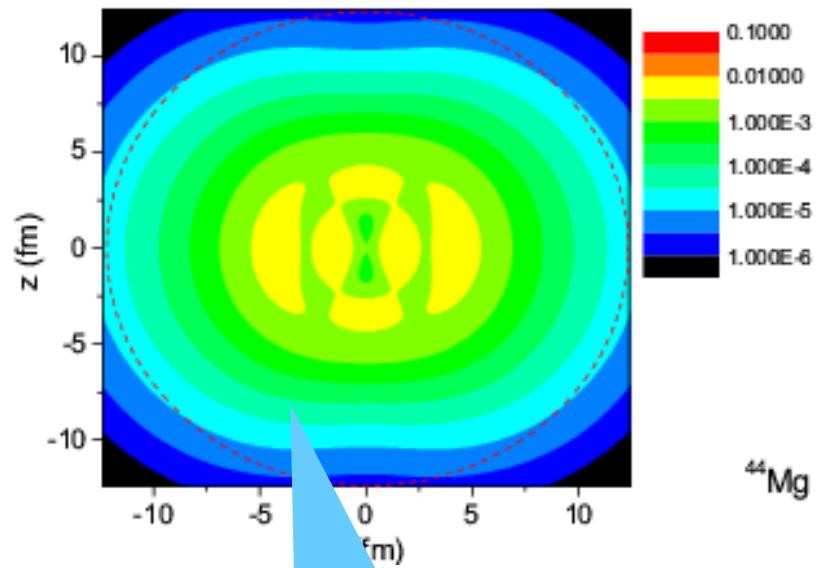


- Prolate deformation
- Large spatial extension in neutron density distribution

# $^{44}\text{Mg}$ : Density of core & halo---shape decoupling



Core: prolate



halo: oblate

# $^{22}\text{C}$ : Puzzles in $S_{2n}$ , $r_m$ & halo configuration

## □ Two-neutron separation energy

Refs.	AME2003	AME2012	Gaudefroy et al. 2012 PRL109-202503	AME2016
$S_{2n}$ (MeV)	$0.420 \pm 0.940\#$	$0.110 \pm 0.060$	$-0.14 \pm 0.46$	$0.035 \pm 0.020$

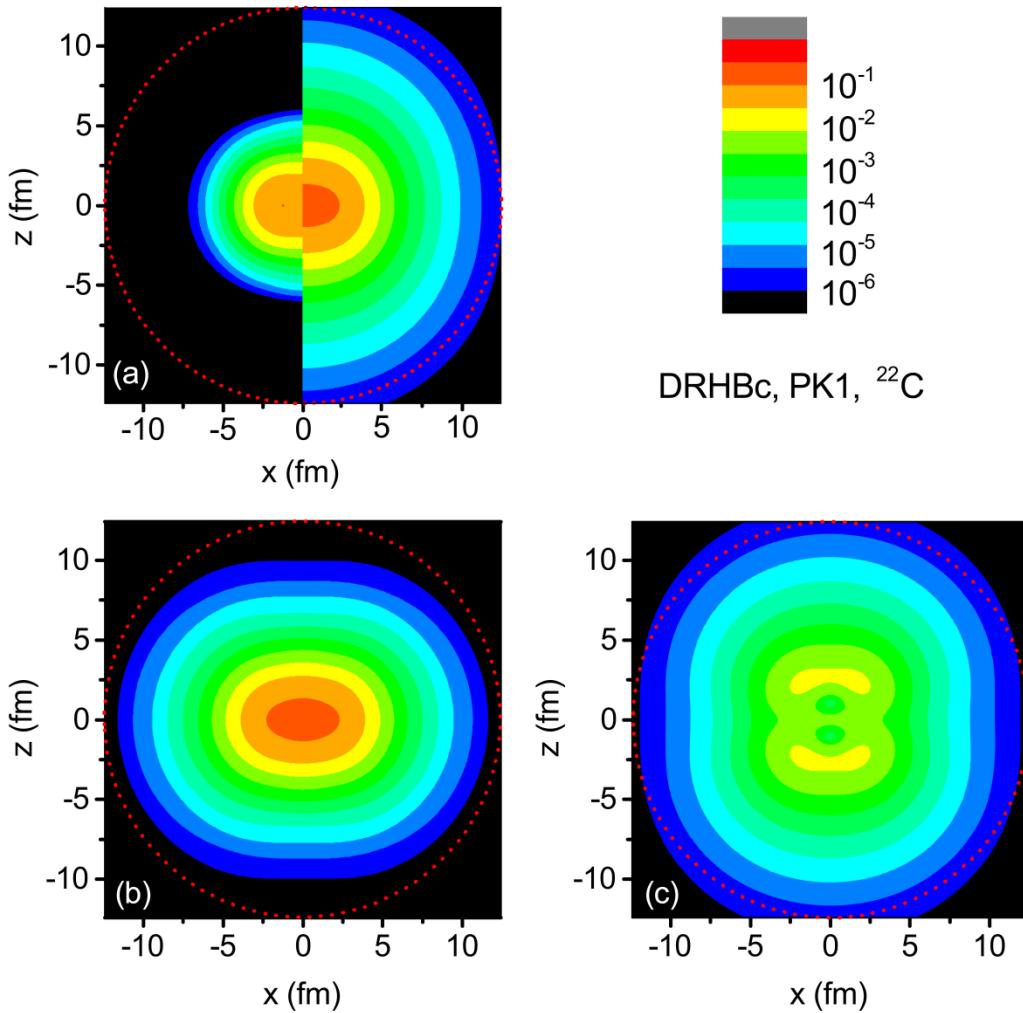
## □ RMS matter radius

Refs.	Tanaka et al. 2010 PRL104-062701	Togano et al. 2016 PLB761-412	Nagahisa&Horiuchi 2018 PRC97-054614	$1.2 A^{1/3} \text{ fm}$
$r_m$ (fm)	$5.4 \pm 0.9$	$3.44 \pm 0.08$	$3.38 \pm 0.10$	$3.36$

## □ Halo configuration

- Inert  $^{20}\text{C}$  w/ 2n in  $2s_{1/2}$  (Horiuchi&Suzuki2006\_PRC74-034311, Ershov et al. 2012\_PRC86-034331, ...)
- Correlated  $^{20}\text{C}$  w/ 2n partly in  $2s_{1/2}$  (Suzuki...2016\_PLB753-199)
- Skyrme Hartree-Fock:  $t_0$  adjusted (Inakura...2014\_PRC89-064316)
- RHFB: no halo (Lu...2013\_PRC87-034311)

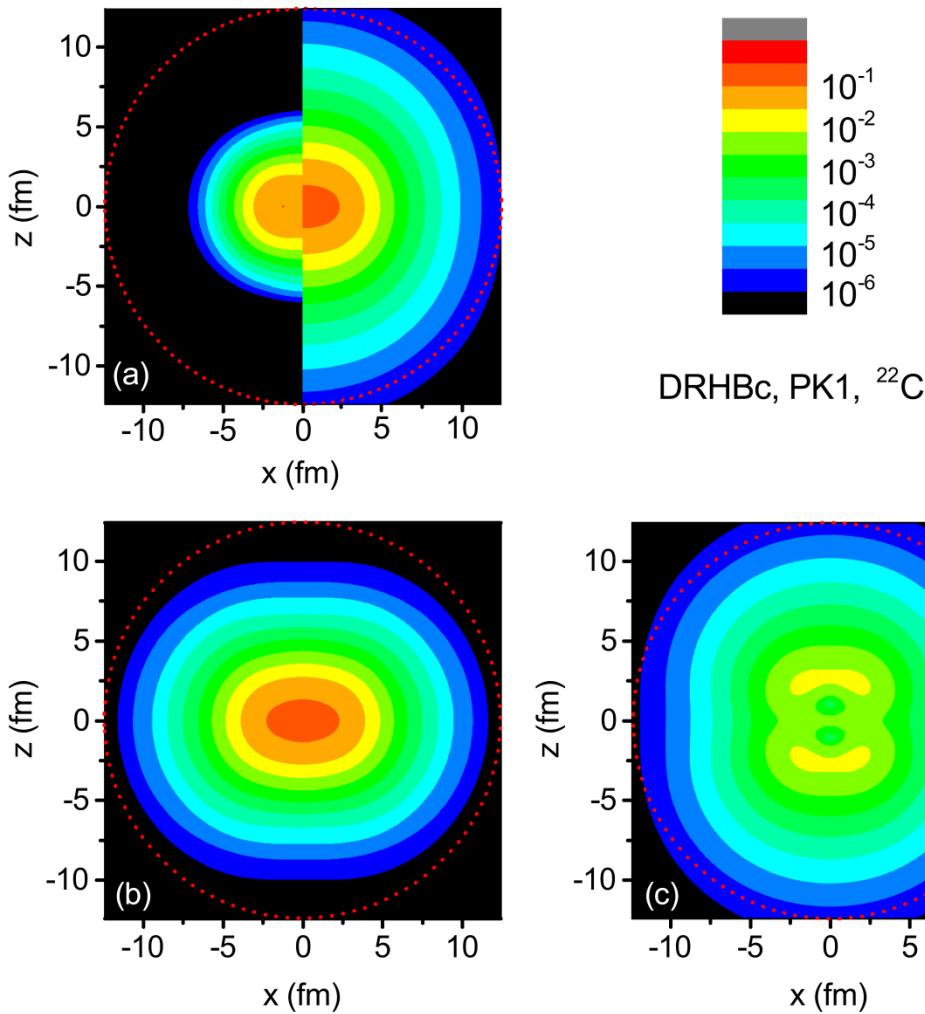
# $^{22}\text{C}$ : Halo (?) & shape decoupling



PK1

$$S_{2n} = 0.43 \text{ MeV}$$
$$r_m = 3.25 \text{ fm}$$
$$\beta_2 = -0.27$$

# $^{22}\text{C}$ : Halo (?) & shape decoupling

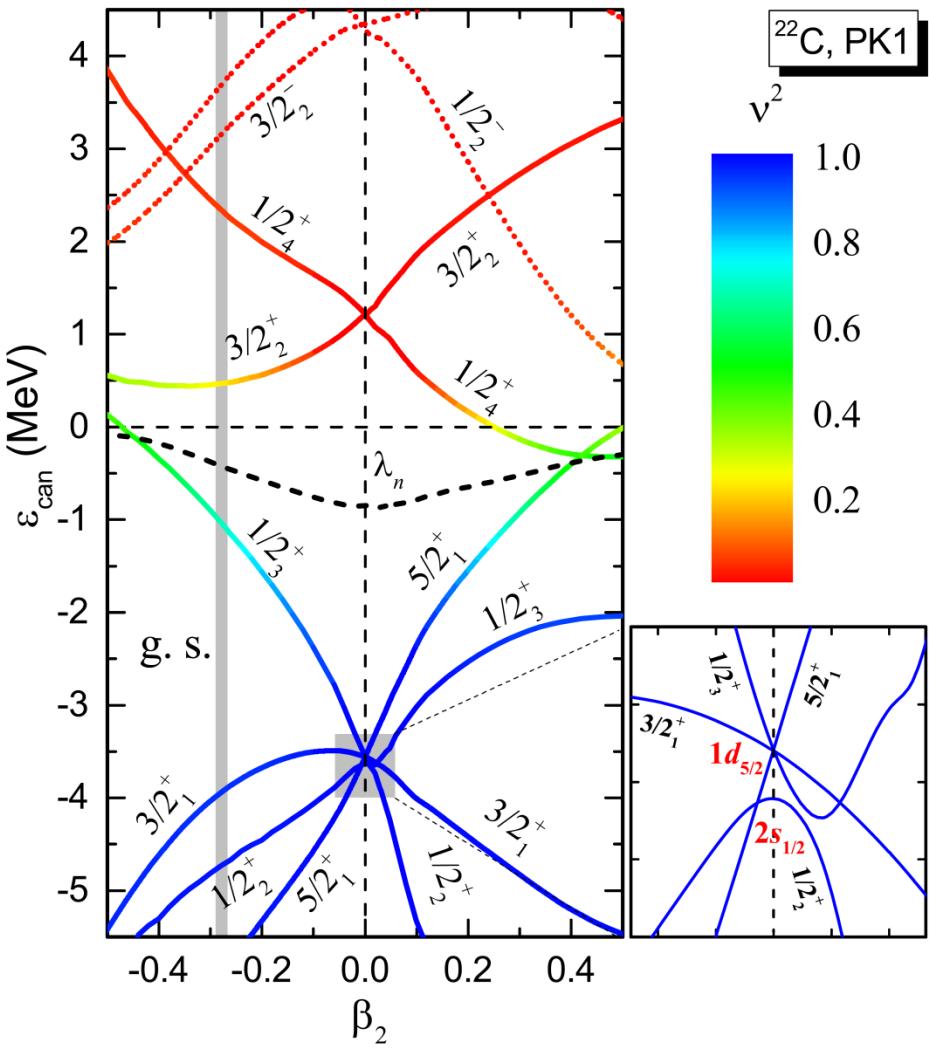


- $2s_{1/2}$ : ~25%  $\Rightarrow$  Halo
- Mixture of ( $2s_{1/2}$ ,  $1d_{5/2}$ )  
 $\Rightarrow$  Prolate halo

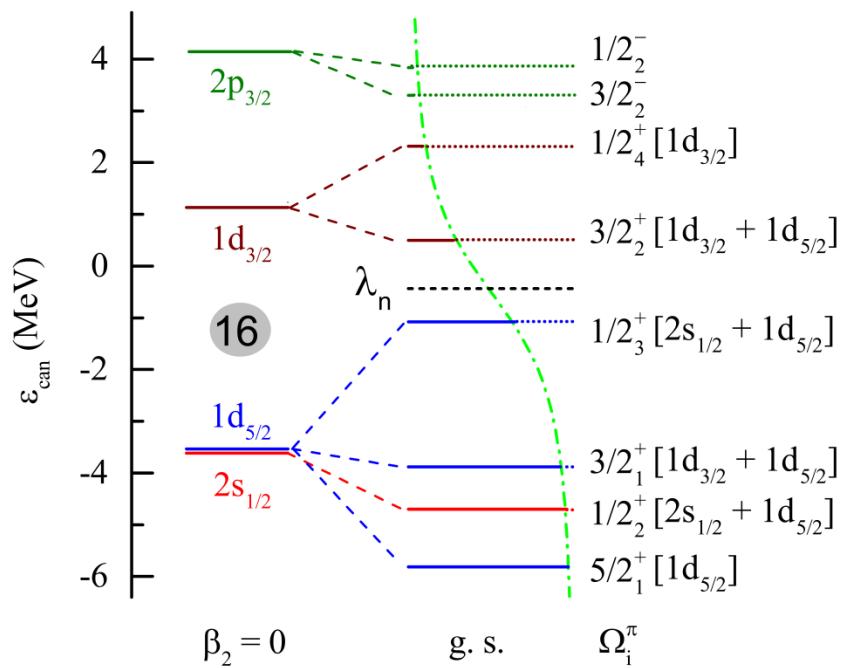
PK1

$$\begin{aligned}S_{2n} &= 0.43 \text{ MeV} \\r_m &= 3.25 \text{ fm} \\\beta_2 &= -0.27\end{aligned}$$

# $^{22}\text{C}$ : Single neutron levels



Inversion of  $(2s_{1/2}, 1d_{5/2})$   
No shell closure at  $N = 16$



# pseudospin symmetry

Volume 30B, number 8

PHYSICS LETTERS

8 December 1969

## PSEUDO LS COUPLING AND PSEUDO $SU_3$ COUPLING SCHEMES \*

A. ARIMA‡

*Department of Physics, State University of New York at Stony Brook, New York 11790, USA*

M. HARVEY

*Theoretical Physics Division, Chalk River Nuclear Laboratories, Ontario, Canada*

and

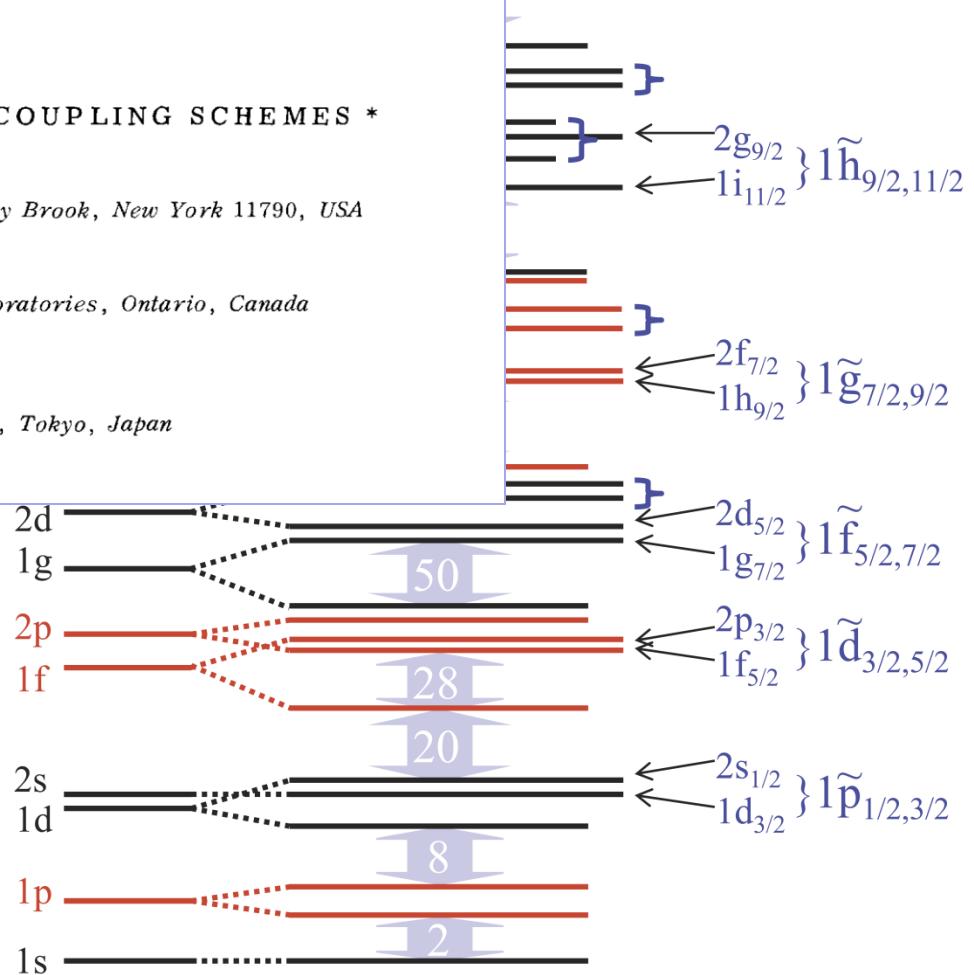
K. SHIMIZU

*Department of Physics, University of Tokyo, Tokyo, Japan*

Received 20 October 1969

Arima\_Harvey\_Shimizu 1969\_PLB30-517

Hecht\_Adler 1969\_NPA137-129



# Justification of PSS in Resonant States

PRL 109, 072501 (2012)

PHYSICAL REVIEW LETTERS

week ending  
17 AUGUST 2012

## Pseudospin Symmetry in Single Particle Resonant States

Bing-Nan Lu (吕炳楠),<sup>1</sup> En-Guang Zhao (赵恩广),<sup>1,2</sup> and Shan-Gui Zhou (周善贵)<sup>1,2,\*</sup>

<sup>1</sup>*State Key Laboratory of Theoretical Physics, Institute of Theoretical Physics, Chinese Academy of Sciences,  
Beijing 100190, China*

<sup>2</sup>*Center of Theoretical Nuclear Physics, National Laboratory of Heavy Ion Accelerator, Lanzhou 730000, China*  
(Received 23 April 2012; revised manuscript received 2 July 2012; published 16 August 2012)

The pseudospin symmetry (PSS) is a relativistic dynamical symmetry connected with the small component of the Dirac spinor. The origin of PSS in single particle bound states in atomic nuclei has been revealed and studied extensively. By examining the zeros of Jost functions corresponding to the small components of Dirac wave functions and phase shifts of continuum states, we show that the PSS in single particle resonant states in nuclei is conserved when the attractive scalar and repulsive vector potentials have the same magnitude but opposite sign. The exact conservation and the breaking of the PSS are illustrated for single particle resonances in spherical square-well and Woods-Saxon potentials.

DOI: 10.1103/PhysRevLett.109.072501

PACS numbers: 21.10.Pc, 03.65.Nk, 21.10.Tg, 24.10.Jv

# Justification of PSS in Renonant States

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**The 8<sup>th</sup> China-Japan Joint Nuclear Physics Symposium  
(CJJNPS 2012)**

**15-19 October 2012 Beijing, China**

# Justification of PSS in Renonant States

The 8<sup>th</sup> China-Japan Joint Nuclear Physics Symposium  
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15-19 October 2012 Beijing, Ch

收件人 主题 日期 大

tanaki Re: an article for Professo... 2012年10月16日 2...

tanaki an article for Professor Ar... 2012年10月16日 2...

an article for Professor Arima

Shan-Gui Zhou

收件人: tanaki

附件: Lu2012\_PRL109-072501.pdf (153 K)

Dear Ms. Tanaki,

Could you please print the attached article and give it to Professor Arima. Thank you very much.

We met yesterday in the 8th China-Japan Joing Nuclear Physics Symposium and he is interested in this article.

With best regards,  
Shan-Gui

---

Prof. Dr. Shan-Gui Zhou  
Institute of Theoretical Physics  
Chinese Academy of Sciences  
55 Zhongguancun Donglu  
Beijing 100190, China  
Tel 86-10-6255 1138 (01)

# pseudospin symmetry

Volume 30B, number 8

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Liang\_Meng\_SGZ 2015  
PhysRep570-1

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Physics Reports 570 (2015) 1–84

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Hidden pseudospin and spin symmetries and their origins in atomic nuclei

Haozhao Liang <sup>a,b</sup>, Jie Meng <sup>a,c,d,\*</sup>, Shan-Gui Zhou <sup>e,f</sup>



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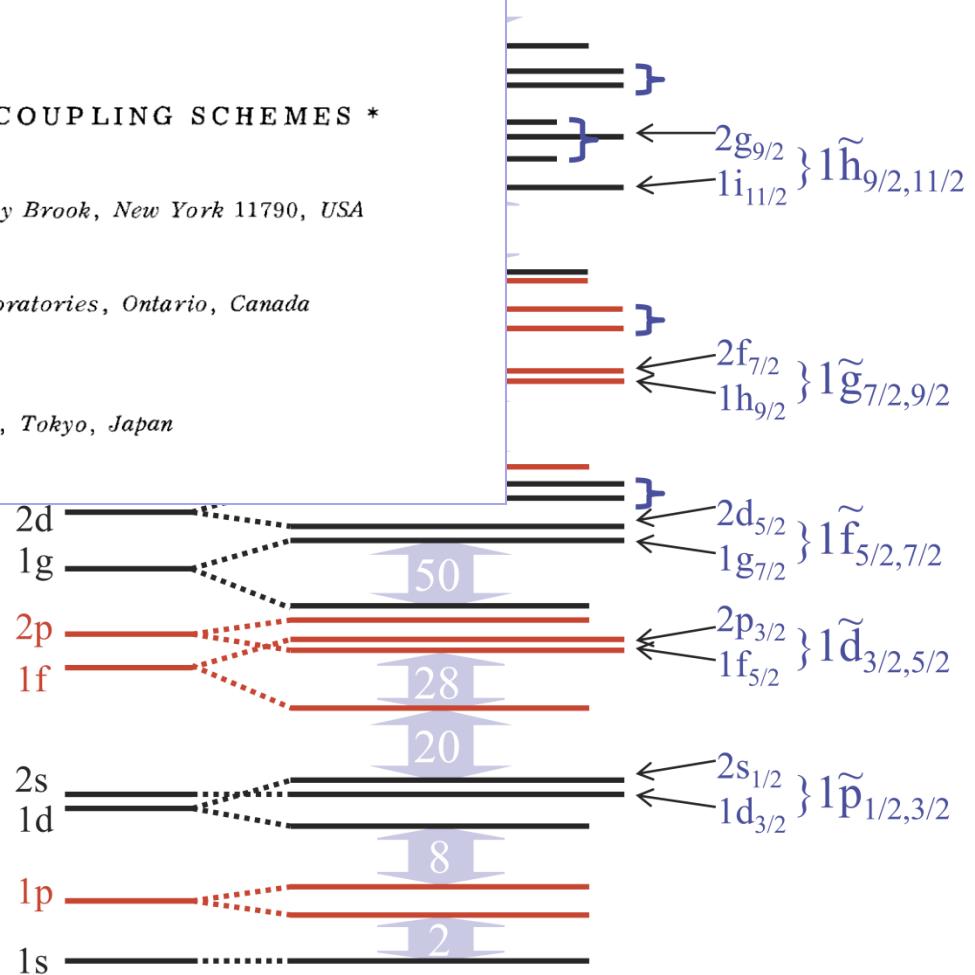
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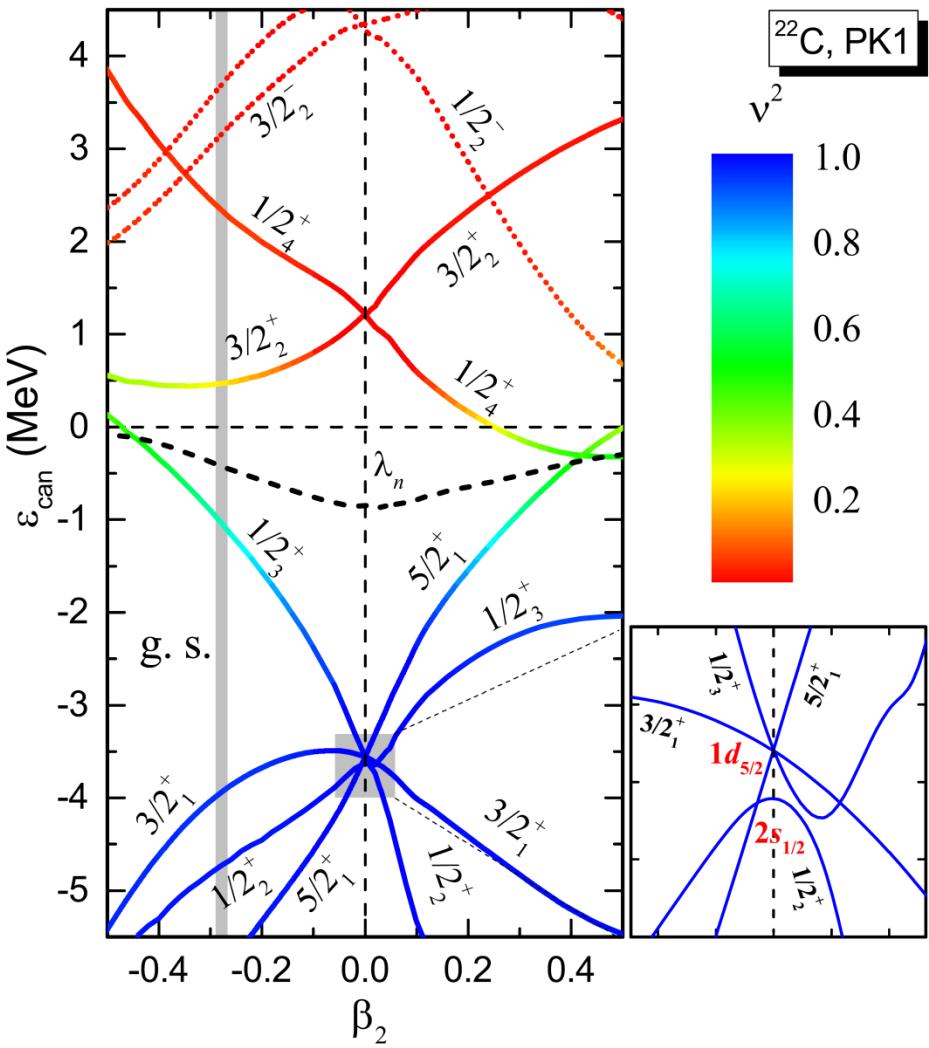
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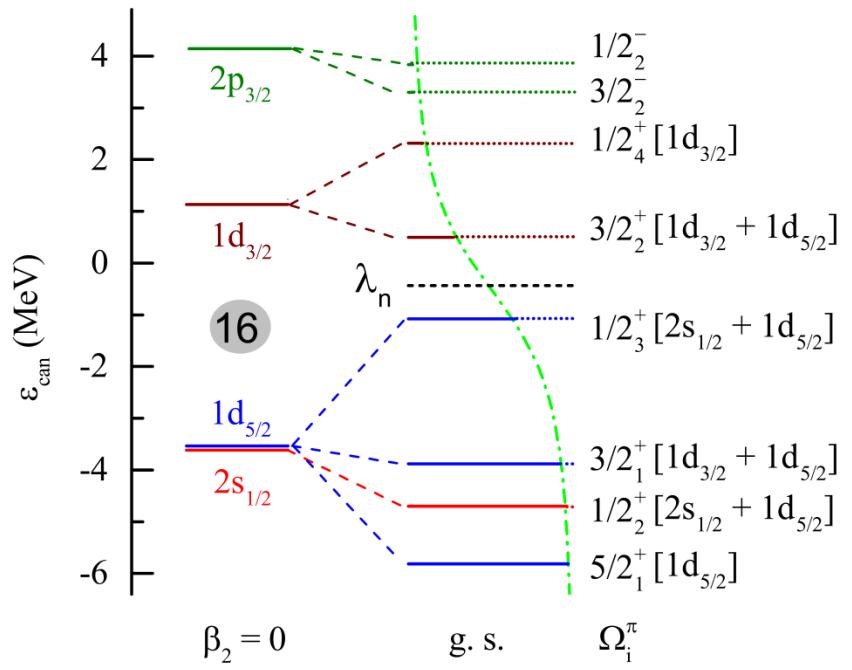
Hecht\_Adler 1969\_NPA137-129



# Breaking of pseudospin symmetry?



Inversion of  $(2s_{1/2}, 1d_{5/2})$   
No shell closure at  $N = 16$



# Breaking of pseudospin symmetry?

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Many symmetries are not made for conservation:  
They can be broken but not conserved.

Liang\_Meng\_SGZ 2015  
PhysRep570-1

Arima\_Harvey\_Shimizu  
1969\_PLB30-517

Hecht\_Adler  
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Physics Reports 570 (2015) 1–84

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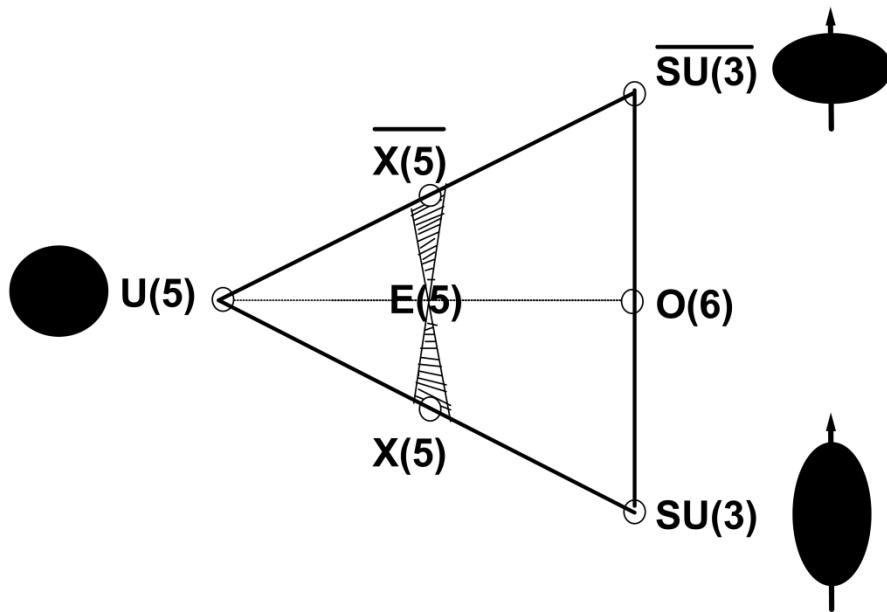


Hidden pseudospin and spin symmetries and their origins in atomic nuclei

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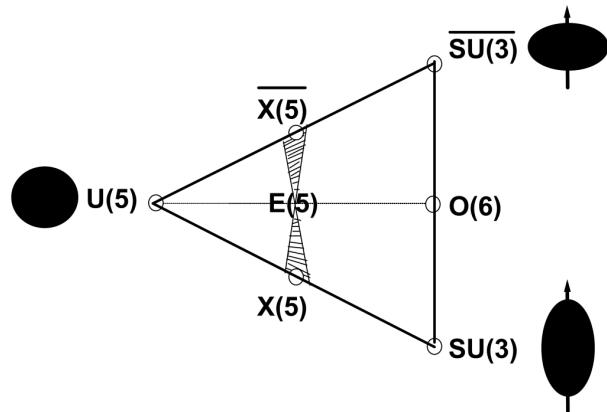
# Extended Casten triangle



# Triangle of Borromean nuclei: $^{11}\text{Li}$ , $^{22}\text{C}$ & $^{44}\text{Mg}$

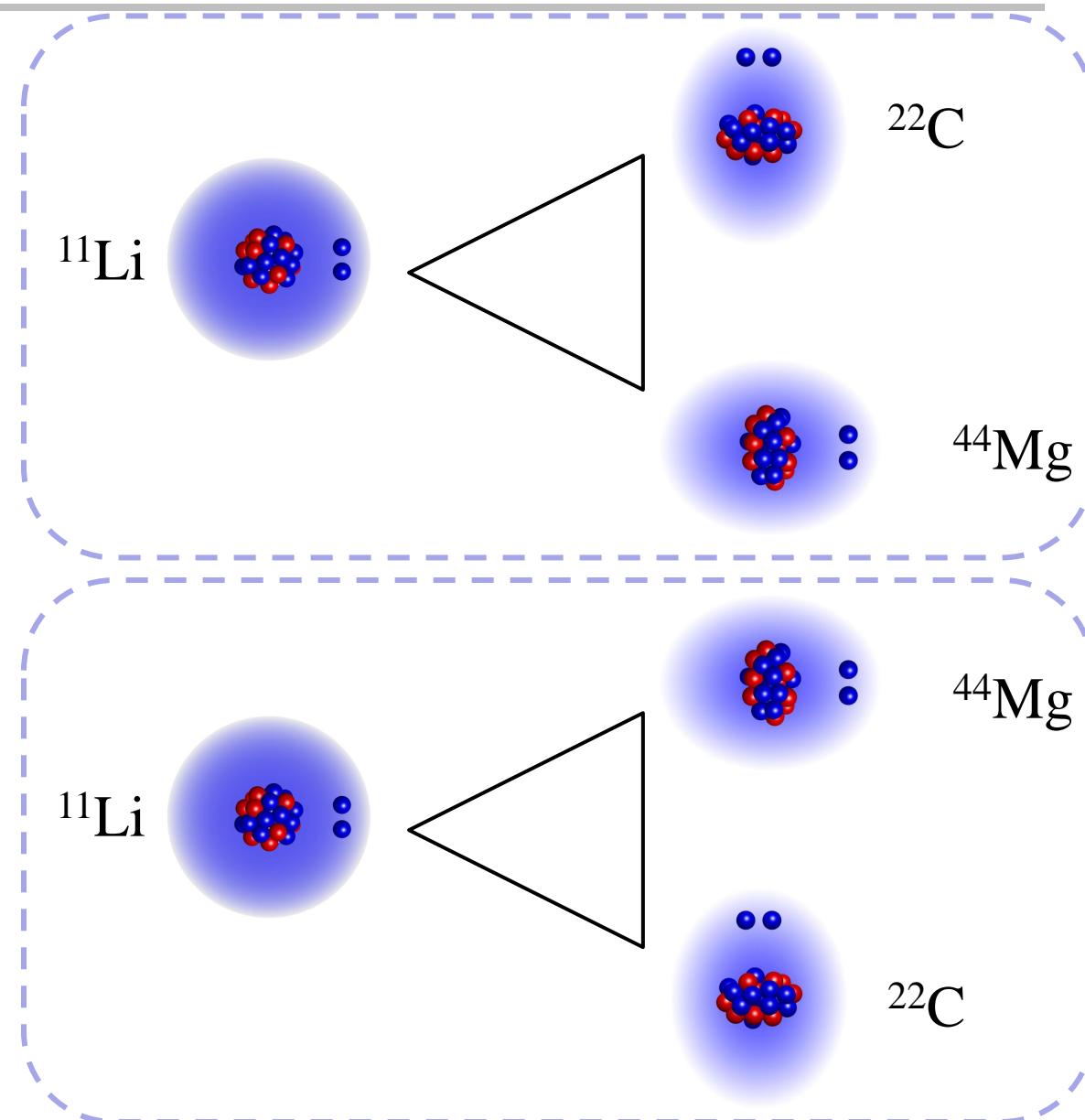
$$^{44}\text{Mg} = ^{22}\text{C} + ^{22}\text{C}$$

$$^{22}\text{C} = ^{11}\text{Li} + ^{11}\text{Li}$$



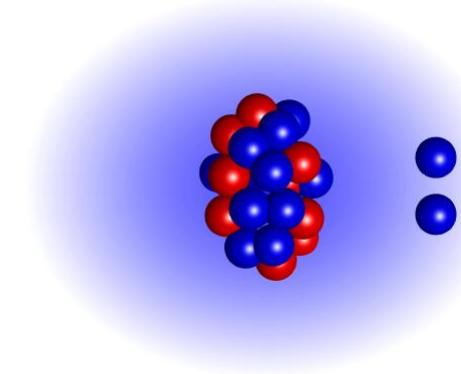
Pan\_Wang\_Huo\_Draayer  
2006\_IJMPE15-1723

Picture(s): courtesy of  
Xin-Hui Wu (吴鑫辉)



# How to probe the shape decoupling?

- Larger cross section
- Narrower momentum distribution
  - Bimodal ?

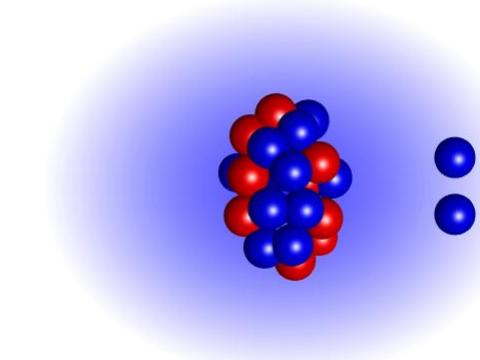


# How to probe the shape decoupling?

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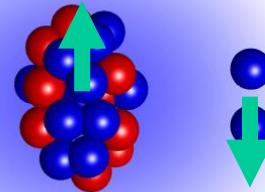
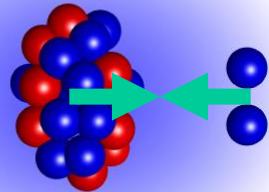
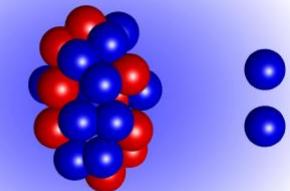
Navin...1997\_PRL81-5089

Sakharuk\_Zelevinsky1998\_PRC61-014609



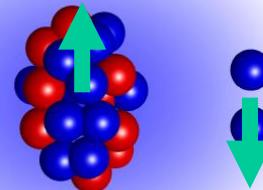
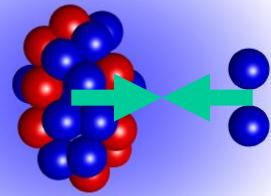
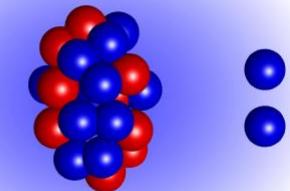
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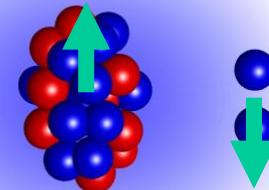
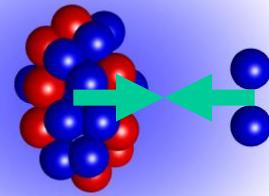
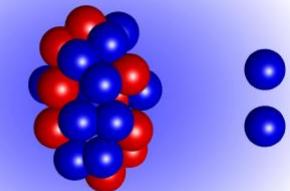
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# Summary & perspectives

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- Deformed relativistic HB theory in a Woods-Saxon basis
  - Occurrence of a halo in deformed nuclei depending on intrinsic structure of valence orbitals
  - $^{44}\text{Mg}$ : prolate core but oblate halo
  - $^{22}\text{C}$ : oblate core but prolate halo
  - $^{11}\text{Li}$ ,  $^{22}\text{C}$  &  $^{44}\text{Mg}$ : triangle of Borromean nuclei ?
- Breaking of pseudospin symmetry ?
- How to probe shape decoupling ?

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祝有马朗人先生  
福如东海，寿比南山！